



Application of Delaunay triangulation to the nearest neighbour method of strain analysis

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Abstract

The nearest neighbour method of strain analysis is re-evaluated and a method for objectively determining nearest neighbours, namely the Delaunay triangulation, is applied. A simulation study and application to a real set of data demonstrates that this approach makes the NNM of strain analysis a practical (and computationally more efficient) alternative to the Fry and associated methods. Once nearest neighbours are selected centre–centre distances can be processed by normalisation and enhancement and the best fit ellipse is determined using a steepest gradient non-linear least squares algorithm applied to the polar equation of a centred ellipse. A simulation study indicates that the technique is a valid one and estimates the strain ellipse well at the 95% confidence interval. Application to a set of natural oolite data shows that there is a systematic variation of error with selection factor and it is suggested that the best estimate of the strain ellipse is obtained by choosing the selection factor which minimises the error.

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1. Introduction

Ramsay (1967, p. 195) introduced the idea that finite strain may be estimated using the distance between centres of adjacent objects, provided the objects were evenly distributed prior to deformation. Implementation of this technique in two dimensions involves determining nearest neighbours, calculating the distance between nearest neighbours (d_i) and the orientation of the line joining them (ϕ_i). Subsequently by plotting d_i against ϕ_i , the minimum and maximum averaged values of d_i (d_{\min} and d_{\max}) could be estimated and divided to provide an estimate of ellipticity R . The orientation of the extensional axis of the strain ellipse is taken to parallel the orientation of d_{\max} . Strain may also be estimated by determining the best-fit least-squares ellipse for the data (Ramsay and Huber, 1983, p. 117; Erslev and Ge, 1990; Hart and Rudman, 1997). Although the nearest neighbour method (NNM) has been shown to work, there are several problems associated with it (Erslev, 1988):

1. No exact definition of what constitutes a nearest neighbour leading to subjectivity in measurement.

2. Difficulties with interpretation.
3. Labour intensity.

As a result there has been minimal application of the Ramsay's NNM.

Graphical methods of strain analysis, which utilise all object–object separations, such as Fry's Method (Fry, 1979), the Normalised Fry Method (Erslev 1988) and the enhanced Normalised Fry Method (Erslev and Ge, 1990) have superseded the NNM due to their simplicity and easy application. In addition, McNaught (1994) extended these methods for application to aggregates of non-elliptical objects. More recently these methods have been incorporated into automated computer based analysis tools including automatic image analysis applications (Aillieres and Champenois, 1994; Aillieres et al., 1995). Although these methods are graphically simple, they are computationally expensive and are of $O(N^2)$, i.e. to process N data takes approximately N^2 steps.

In this paper the NNM is revisited in the context of the relatively new subject of computational geometry. The object of this paper is to demonstrate that the NNM readily lends itself to analysis by the methods of computational geometry such as Delaunay triangulation and that the finite strain may be determined by statistical analysis of the resulting data. In order to distinguish between Ramsay's

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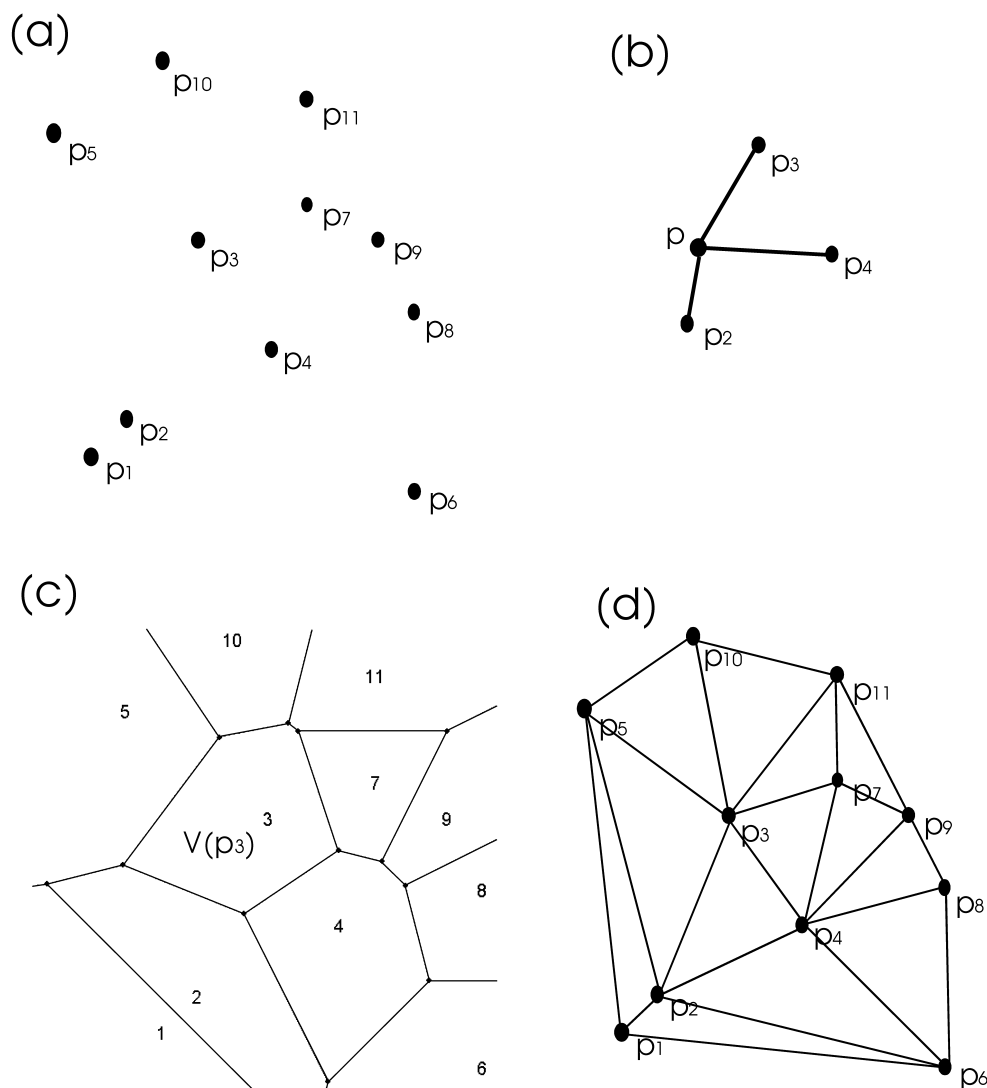


Fig. 1. (a) An example set of 10 points $P = \{p_1, p_2, \dots, p_{10}\}$ in the plane. (b) When creating a Voronoi diagram the plane is partitioned by assigning every point in the plane to its nearest neighbour. In this subset of points the arbitrary point p is nearest to p_2 . In constructing a Voronoi diagram every point in the plane is assigned to its nearest neighbour, points with more than one nearest neighbour form part of the Voronoi diagram. (c) The Voronoi diagram corresponding to the example set of points. (d) The Delaunay triangulation of the example set of points.

NNM and that developed here, the current method will be referred to as the DTNNM (standing for Delaunay triangulation nearest neighbour method). Use of the Delaunay triangulation and statistical analysis removes subjectivity in measurement and interpretation. Furthermore, the DTNNM is computationally more efficient by comparison with the various Fry methods (Fry, 1979; Erslev, 1988; Erslev and Ge, 1990; McNaught, 1994). This is because calculating a Delaunay triangulation is $O(MnN)$ in the worst case (Perparata and Shamos, 1985, p. 123; Edelsbrunner, 1987, p. 305; O' Rourke, 1993, p. 200; Huang and Shih, 1998), i.e. it is 50 times faster for 100 data and 300 times faster for 1000 data. In addition it also automatically selects a subset of nearest neighbours. The difference between an $O(N^2)$ and an $O(MnN)$ algorithm is not an issue on most modern personal computers for a typical geological dataset; however, it may be an issue for large datasets or

collections of datasets particularly in the light of the development of automated image analysis techniques (Aillieres and Champenois, 1994; Aillieres et al., 1995).

2. Delaunay triangulation, Voronoi diagram and convex hull

A Delaunay triangulation is the set of lines joining a set of points together such that each point is joined to its nearest neighbours (Perparata and Shamos, 1985, p. 209; O' Rourke, 1993, p. 175). Before elaborating further on Delaunay triangulations, it is instructive to consider a closely related property of a set of points, namely the Voronoi diagram. Consider a set of n nodes $P = \{p_1, p_2, p_3, \dots, p_n\}$ (see Fig. 1a for an example set) and that the plane is partitioned by assigning every point in it to its nearest

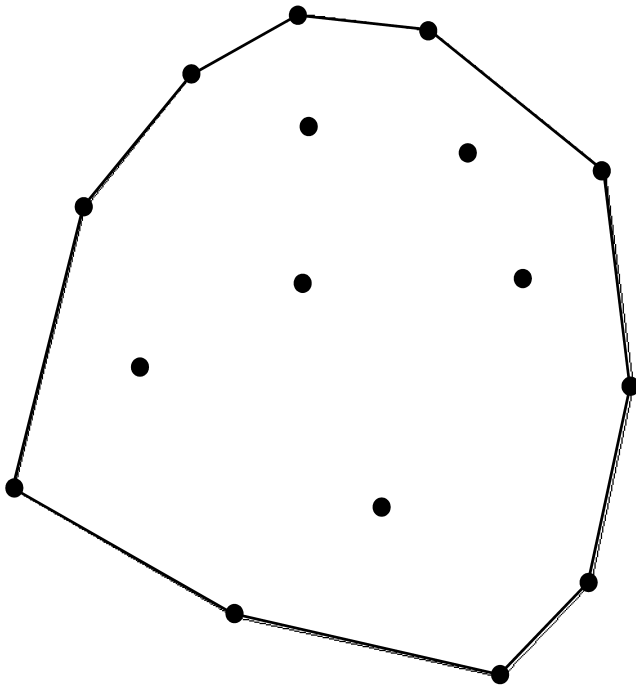


Fig. 2. An example of the convex hull of a set of points. Notice that the hull is everywhere convex and fully encloses the points in the set.

neighbour in P , for example in Fig. 1b the point p is assigned to the node p_2 , its nearest neighbour. All points whose nearest neighbour is p_i form a region of the plane termed a Voronoi region denoted by $V(p_i)$ (see Fig. 1c). Certain points do not have a unique neighbour and may be equally close to two or more nodes, all such points form part of the Voronoi diagram and lie along the lines in Fig. 1c. A Voronoi diagram consists of a set of adjacent straight-sided polygons that divide the plane into regions of closeness to a given node. Voronoi diagrams have found many applications in diverse subject areas including crystallography where Voronoi diagrams are used to simulate constant crystal growth from dispersed sites of nucleation.

Once a Voronoi diagram has been devised, a Delaunay triangulation may be superimposed by joining together nodes, whose Voronoi regions touch, by a line (Fig. 1d). The line joining the central nodes of two adjacent and touching Voronoi regions is the local shortest path between the two nodes, because each Voronoi region consists of the points which are closest to the central node. Many algorithms exist which compute the Delaunay triangulation (Huang and Shih, 1998) and Voronoi diagram of a set of nodes and the interested reader is referred to O' Rourke (1993) for a thorough treatment.

A set of lines connecting nearest neighbours is the result of computing a Delaunay triangulation. This is exactly the information required to apply Ramsay's NNM of strain analysis. Because a Delaunay triangulation is calculated computationally, it eliminates the time consuming and subjective work that was previously associated with the Ramsay's NNM.

The convex hull of a set of points in two dimensions is related to the above notions and is defined as the enclosing convex polygon with the smallest perimeter and area (O' Rourke, 1993, p. 70; see Fig. 2). The Voronoi diagram, Delaunay triangulation and convex hull all have three-dimensional counterparts so that the methods introduced here for two dimensions can be quite easily applied to higher dimensional problems.

3. Methodology

Objects used in the analysis are idealised as ellipses mainly for computational convenience in the simulation study and because the natural example studied consists of elliptical objects. The primary prerequisite for both the Fry method and the DTNNM is a set of points that were anti-clustered prior to deformation and, equally important, that the same set of points can be identified in the deformed state. Fry (1979) chose to use object centroids for this purpose, which are readily calculated for arbitrary shapes. In general, it is not necessary to first approximate objects of arbitrary shape by ellipses prior to application of the Fry method or DTNNM (as suggested by Erslev and Ge (1990)) because the centroid of an arbitrary shape is as good (if not better) a point as the centre of the fitted ellipse. Moreover, although it may be computationally convenient to calculate the enhancement factor (Erslev and Ge, 1990) for the purpose of simulation using ellipse radii as described below, these radii are readily calculable for arbitrary shapes also. Therefore like the Fry method, the DTNNM applies to objects of arbitrary shape.

The raw data required for the analysis is $(a_i, b_i, x_i, y_i, \phi_i)$, where a_i and b_i are the half-lengths of the long and short axes of each ellipse, respectively, x_i and y_i are the coordinates of the centre of each ellipse, ϕ_i is the orientation of the long axis and i takes values between 1 and n , where n is the total number of objects measured. Once the data is acquired the following processing is applied.

Using centre co-ordinates (x_i, y_i) , the Delaunay triangulation is calculated resulting in a data set $D = ((x_i, y_i), (x_j, y_j))$. Both i and j take values between 1 and n and $i \neq j$ (i.e. a centre cannot be connected to itself). D represents the set of connected nearest neighbours. As illustrated below with an example (Fig. 4a), the Delaunay triangulation does not always join nearest neighbours near the edge of the data set. These rogue values may be eliminated by subtracting points which form part of the convex hull from the Delaunay triangulation (see Fig. 4b) and gives a modified Delaunay triangulation (D_m).

D_m is processed further by applying the enhancement of Erslev and Ge (1990). The length (d_{ij}) and orientation (θ_{ij}) of the line segment joining each pair of points (x_i, y_i) and (x_j, y_j) are calculated. The length of the elliptical radius along θ_{ij} for each ellipse (r_i and r_j) is calculated using the polar

equation of an ellipse centred on the origin, i.e.:

$$e_i = \sqrt{1 - \frac{b_i^2}{a_i^2}} \quad (1)$$

$$r_i = a_i \sqrt{\frac{1 - e_i^2}{1 - e_i^2 \cos^2(\theta_{ij} - \phi_i)}} \quad (2)$$

Note that if the objects in question were not well approximated by an ellipse, these distances could also be calculated by measurement. The enhancement factor (ef_{ij}) for each pair of points is given by (Erslev and Ge, 1990):

$$ef_{ij} = \frac{d_{ij}}{r_i + r_j} \quad (3)$$

and is a measure of the adjacency of two ellipses. For example, $ef_{ij} = 1$ implies that the two ellipses are touching along the θ_{ij} direction. The modified Delaunay data set (Dm) is reduced by only accepting points such that $ef_{ij} \leq sf$, where sf is the selection factor of Erslev and Ge (1990). Enhancing the data in this way removes data relating to objects which are far apart.

The remaining data are normalised according to the method of Erslev (1988). The average radius of each ellipse is calculated using $\bar{r}_i = \sqrt{a_i b_i}$ and the normalised distance between two ellipse centres (dn_{ij}) is given by:

$$dn_{ij} = \frac{d_{ij}}{\bar{r}_i + \bar{r}_j} \quad (4)$$

The final data set consists of a set of points (dn_{ij}, θ_{ij}) interpreted as a polar set of data, i.e. length and orientation, centred on the origin. Each point in the final data set actually represents two points in so far as the data are non-directional, i.e. both (dn_{ij}, θ_{ij}) and ($-dn_{ij}, \theta_{ij}$) are valid data points.

To estimate the axial ratio of the strain ellipse (R_s), the best fit ellipse is calculated for ($\pm dn_{ij}, \theta_{ij}$). Other authors (Erslev and Ge, 1990; Hart and Rudman, 1997) have utilised linear least squares for fitting elliptical data to the Cartesian equation of an ellipse centred on the origin. However, due to the fundamentally polar nature of the ellipse (and also the data obtained from the DTNNM), a method based on the polar equation of an ellipse is advocated here. In addition, because the equation of an ellipse in any co-ordinate system is essentially non-linear, a steepest gradient non-linear least squares algorithm (Draper and Smith, 1998, p. 513; Mulchrone, 2001) is used to obtain the best fit ellipse.

The equation for an ellipse centred on the origin with semi-major axis a , semi-minor axis b and whose long axis makes an angle ϕ with the x -axis (counter-clockwise taken as positive) can be written in polar co-ordinates as:

$$r = a_i \sqrt{\frac{1 - e^2}{1 - e^2 \cos^2(\theta - \phi)}} \quad (5)$$

where r is the radial length along an orientation θ and e is the

eccentricity defined as:

$$e = \sqrt{1 - \frac{b^2}{a^2}} \quad (6)$$

Application of trigonometric identities and simple rearrangement allows Eq. (5) to be rewritten as:

$$\frac{1}{r^2} = \left(\frac{1}{b^2} - \frac{e^2}{2b^2} \right) - \frac{e^2}{2b^2} \cos(2\phi) \cos(2\theta) - \frac{e^2}{2b^2} \sin(2\phi) \sin(2\theta) \quad (7)$$

However, this may be idealised as follows:

$$\frac{1}{r^2} = C - A \cos(2\theta) - B \sin(2\theta) \quad (8)$$

where

$$A = \frac{e^2}{2b^2} \cos(2\phi),$$

$$B = \frac{e^2}{2b^2} \sin(2\phi),$$

and

$$C = \left(\frac{1}{b^2} - \frac{e^2}{2b^2} \right).$$

So, given a set of data ($1/r_i^2, \theta_i$), a non-linear least squares multiple regression can be applied to estimate the parameters A , B and C . The axial ratio (R_s) and orientation of the long axis (ϕ_s) of the strain ellipse may be calculated as follows:

$$\tan(2\phi_s) = \frac{B}{A} \quad (9)$$

$$R_s = \sqrt{\frac{C \cos(2\phi_s) + A}{C \cos(2\phi_s) - A}} \quad (10)$$

note that depending on whether A is positive or negative R_s may need to be inverted to get a value greater than one. As usual all angles should be in radians.

In terms of the data resulting from the DTNNM the data set (dn_{ij}, θ_{ij}) is transformed to ($1/dn_{ij}^2, \theta_{ij}$) and then fitted to Eq. (8). It is noted that a linear multiple least squares regression could also be applied to Eq. (8) by letting $X = \cos(2\theta)$ and $Y = \sin(2\theta)$ so that the data set (dn_{ij}, θ_{ij}) is transformed to ($1/dn_{ij}^2, X_{ij}, Y_{ij}$). This approach produces equivalent results.

4. Simulation study

4.1. Introduction

This study was motivated by a desire to verify and examine the properties of the DTNNM developed here. For

Table 1
List of symbols used in the paper and their meaning

Parameter	Description
R_{sact}	Actual strain ratio
sf	Selection factor
n	Number of objects used in the analysis
N_{del}	Number of centre–centre pairs after modified Delaunay triangulation
N_{proc}	Number of centre–centre pairs after enhancement
R_{scalel}	Lower 95% confidence interval for calculated R_s
R_{scale}	Calculated R_s
R_{scaleu}	Upper 95% confidence interval for calculated R_s
ϕ_{scalel}	Lower 95% confidence interval for calculated ϕ_s
ϕ_{scale}	Calculated ϕ_s
ϕ_{scaleu}	Upper 95% confidence interval for calculated ϕ_s

each simulation an initial data set was generated using a simple sequential inhibition process (see below and Diggle, 1983, p. 60). Six initial datasets were generated consisting of $n = 175, 205, 240, 270, 295$ and 330 circles with radii randomly selected from the $[5,10]$ interval and centres lying within a square region ranging from $(0,0)$ to $(300,300)$. Each dataset was then progressively strained in 0.2 steps up to a maximum of $R_s = 5.8$, such that the long axis of the strain ellipse was parallel to the x -axis i.e. 0° . At each strain step, the data were analysed as described above and the best-fit strain ellipse calculated for $sf = 1.01$ to 1.3 in 0.01 steps. A steepest gradient non-linear least squares algorithm pro-

vided with Mathematica (i.e. NonlinearRegress; see Abell et al., 1999, pp. 528–530) was used to calculate the best fit strain ellipse and also gave 95% confidence intervals for the calculated values of R_s and ϕ_s . For each value of R_s and sf the parameters in Table 1 were recorded.

4.2. Generating anti-clustered data

Fundamental to the DTNNM and the method of Fry (1979) is the constraint that input data should be anti-clustered, meaning that in some sense the centres of objects are evenly distributed throughout the plane. In statistical terms all point distributions are realisations of spatial point processes (Diggle, 1983, pp. 46–69); however, the initial anti-clustered distribution required for strain analysis is also stationary and isotropic. The initial anti-clustered point distribution is stationary in so far as the calculated strain value is expected to be invariant under an arbitrary translation and it is isotropic because it is invariant for arbitrary rotation. However, after strain the point distribution is anisotropic, because points are closer to their neighbours along the shortening direction and further apart along the stretching direction.

The required initial anti-clustered point distribution is a realisation of a simple sequential inhibition (SSI) process (Diggle, 1983, p. 60). Generating a set of anti-clustered points using the SSI process proceeds as follows:

1. Suppose we wish to generate a point distribution in some convex plane area A , e.g. a rectangle or circle.
2. Let us assume that the points to be generated are centres, $p_i = (x_i, y_i)$, of some objects s_i and that we wish to generate n points.
3. We proceed by generating an object s_i with centre p_i by randomly generating the co-ordinates (x_i, y_i) . If required we can also randomly generate the parameters defining the shape and orientation of the object, e.g. in the case of a circle the radius might be randomly generated.
4. Before accepting the newly created point it is necessary to check that the new object s_i does not intersect any of the existing objects (i.e. s_1 to s_{i-1}). If the object does not

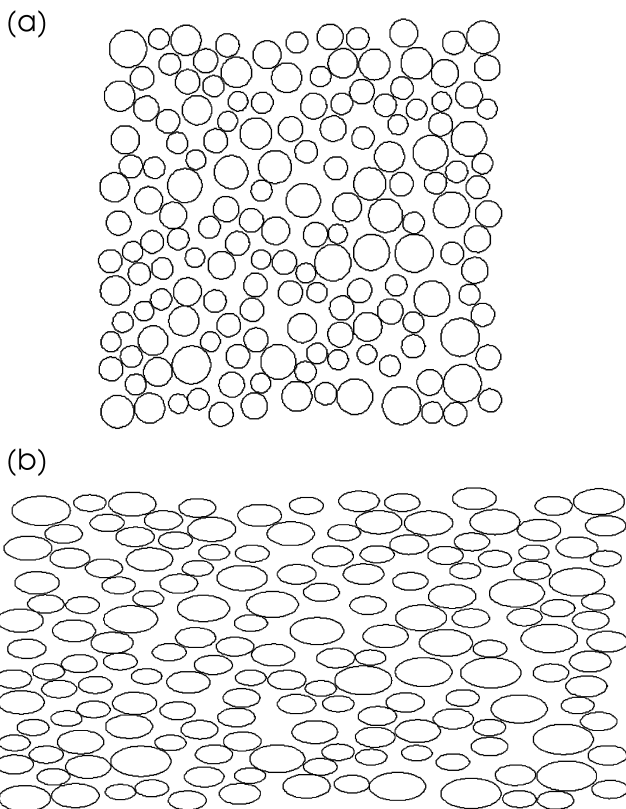


Fig. 3. (a) A set of anti-clustered circular data generated using the SSI process described in the text. (b) The same set of data with a strain of $R_s = 2.0$.

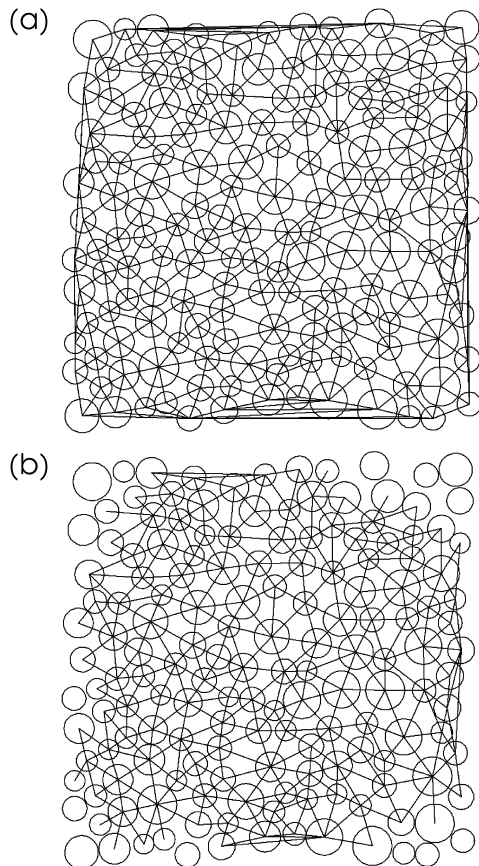


Fig. 4. (a) The data of Fig. 3a along with the calculated Delaunay triangulation. (b) The data of Fig. 3b along with the modified Delaunay triangulation whereby lines connected to points on the convex hull are omitted.

intersect then it is accepted otherwise a new object is generated until such time as it is accepted.

5. Continue to generate objects until such time as the required number of points are acquired.

Clearly there is a limit to the number of objects that can fit inside the finite area A and any implementation of the algorithm should include an opt out clause based on the number of attempts taken to generate an object. In order to test the validity of the DTNNM of strain analysis, initial anti-clustered data sets were generated using the SSI process for circle objects.

The method described above for generating anti-clustered point distributions is similar to that of Fry (1979) and Crespi (1986). Furthermore it is noted that there are alternatives to the SSI process described above such as generating 3D anti-clustered objects and taking a 2D section (Erslev, 1988).

4.3. Example simulation

Fig. 3a shows an example distribution generated with the SSI process. In this example 160 circles, with radius randomly chosen between 5 and 10, were generated to fit

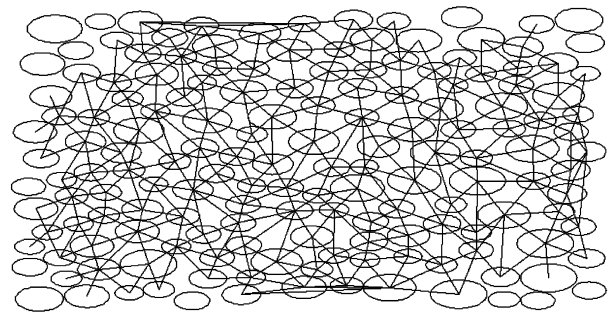


Fig. 5. The strained distribution of Fig. 3b with the corresponding modified Delaunay triangulation.

within a 200×200 square. The distribution after being strained to an axial ratio of two is illustrated in Fig. 3b. Next the Delaunay triangulation is calculated for the unstrained distribution (Fig. 4a). It is readily apparent that in the centre nearest neighbours are joined; however, towards the edges there are points joined which are not nearest neighbours. This is due to the shape of the area used. It may be suggested that a circular or elliptical area would give better results; however, this is a rather impractical restriction. Instead, these edge effects are negated by calculating the convex hull of the data and excluding any line containing a point on the convex hull. The resulting modified Delaunay triangulation is illustrated in Fig. 4b and although there are some rogue connections remaining, their number is severely reduced. A modified Delaunay triangulation was generated for the strained distribution in Fig. 3b and is shown in Fig. 5. It is emphasised that this triangulation was calculated independently of the unstrained triangulation (i.e. this is not a strained triangulation, it is a triangulation of the strained data). Notice that the strained triangulation by and large connects the same points together. There are differences and in general the number of lines in the triangulation decreases with increasing imposed strain. This demonstrates that the Delaunay triangulation is not invariant under strain.

The variance of the Delaunay triangulation with increasing strain is due to a problem recognised by Fry (1979). During homogeneous strain, points which were initially nearest neighbours are pushed further apart along the extensional strain axis but are brought closer together along the shortening direction. Therefore as strain increases, nearest neighbours, and also the Delaunay triangulation, will change. On this basis it is expected that the precision of the DTNNM should deteriorate with increasing strain. However, variance of the Delaunay triangulation is not the only source of error and there will also be random errors due to sampling and deviations from ideal assumed behaviour, e.g. homogeneity and interaction of strain markers.

In order to calculate the strain associated with a particular point distribution, the set of lines connecting nearest neighbours as calculated using the modified Delaunay triangulation are plotted such that each point of each line is translated to the origin and the other point is included in a new transformed data set. Fig. 6a

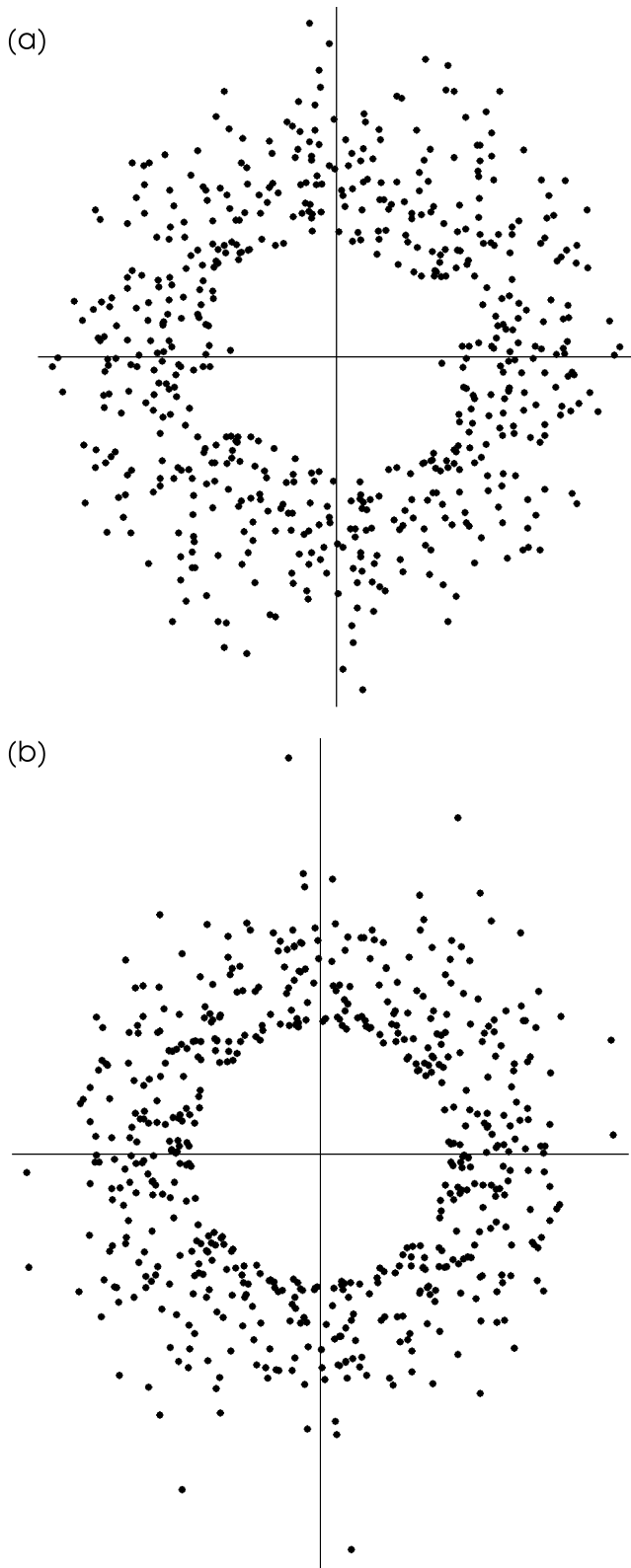


Fig. 6. (a) Plot of the data from a modified Delaunay triangulation corresponding to the unstrained example in Fig. 4b. (b) The same data after normalisation.

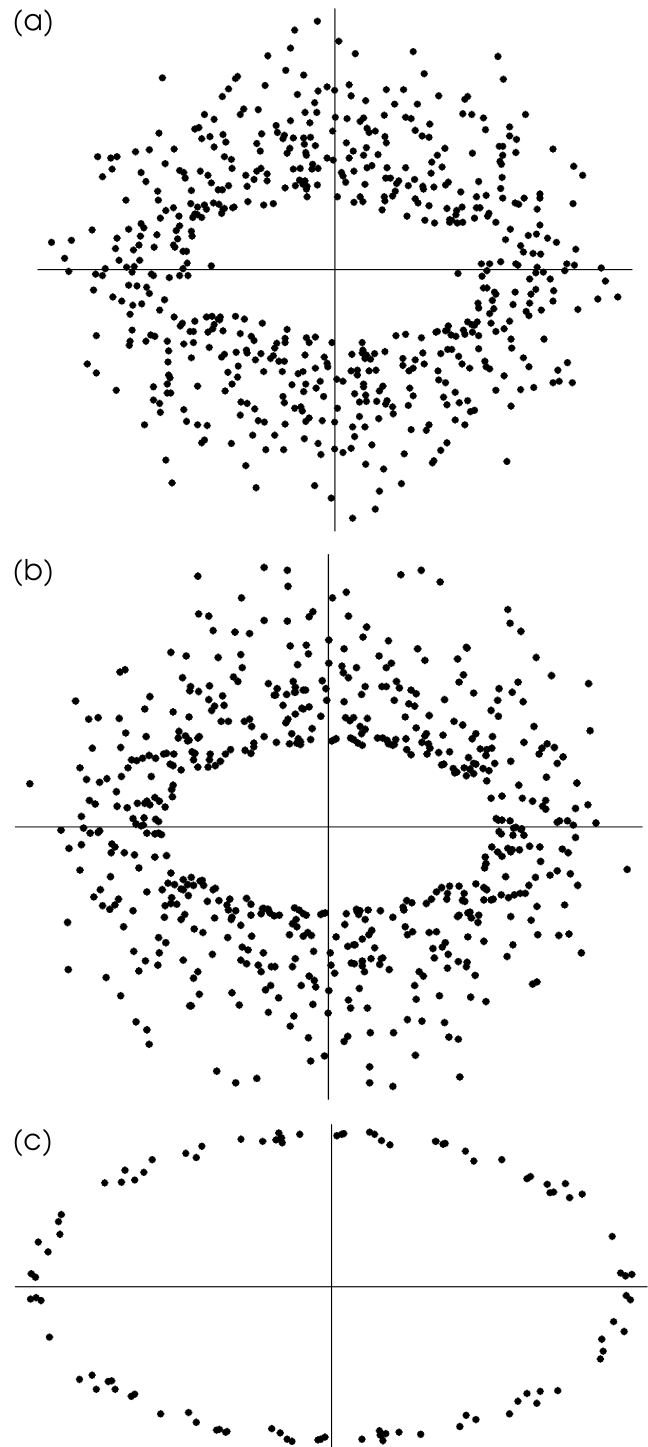


Fig. 7. (a) Plot of the data from a modified Delaunay triangulation corresponding to the strained example in Fig. 5. (b) The same data after normalisation. (c) The same data after enhancement with $sf = 1.1$.

illustrates this plot for the unstrained data and it clearly exhibits a circular pattern. In Fig. 6b the data is normalised according to Erslev (1988) and the central void is very evidently a circle, as expected. Plots for the strained data are shown in Fig. 7. The raw data is shown in Fig. 7a and the normalised data are shown Fig. 7b.

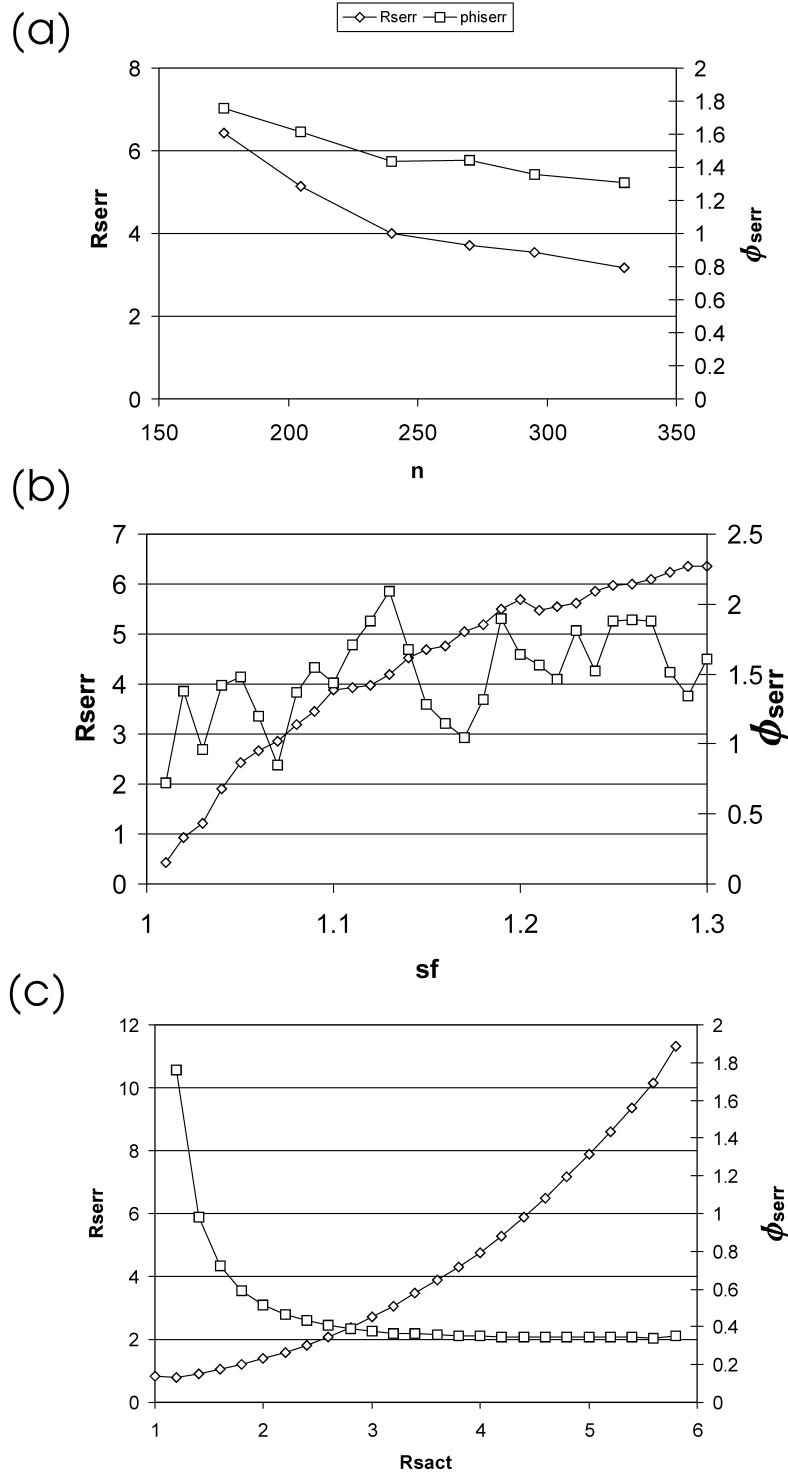


Fig. 8. (a) Variation of average error (R_{serr} and ϕ_{serr}) with n . (b) Variation of average error (R_{serr} and ϕ_{serr}) with increasing selection factor, sf . (c) Variation of average error (R_{serr} and ϕ_{serr}) with increasing imposed strain (R_{sact}).

However, the best result is obtained by applying the enhancement suggest by [Erslev and Ge \(1990\)](#) whereby only objects in close proximity are chosen for analysis. The enhanced normalised plot is shown in [Fig. 7c](#) for a selection factor of 1.1 and the data closely follow an ellipse with axial ratio of two, i.e. the strain ellipse.

4.4. Results

Because a large volume of data was generated by the simulation study, it is impractical to report it in a detailed graphical format. Therefore the average error in R_{scalc} and ϕ_{scalc} was calculated for n , sf and R_{sact} , where the average

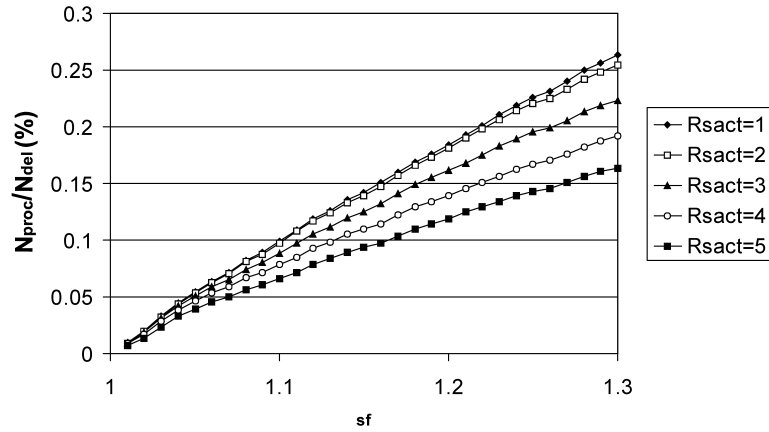


Fig. 9. The percentage ratio of the number of data after processing (N_{proc}) to the number of data from the modified Delaunay triangulation (N_{det}) as a function of both sf and imposed strain (R_{sact}).

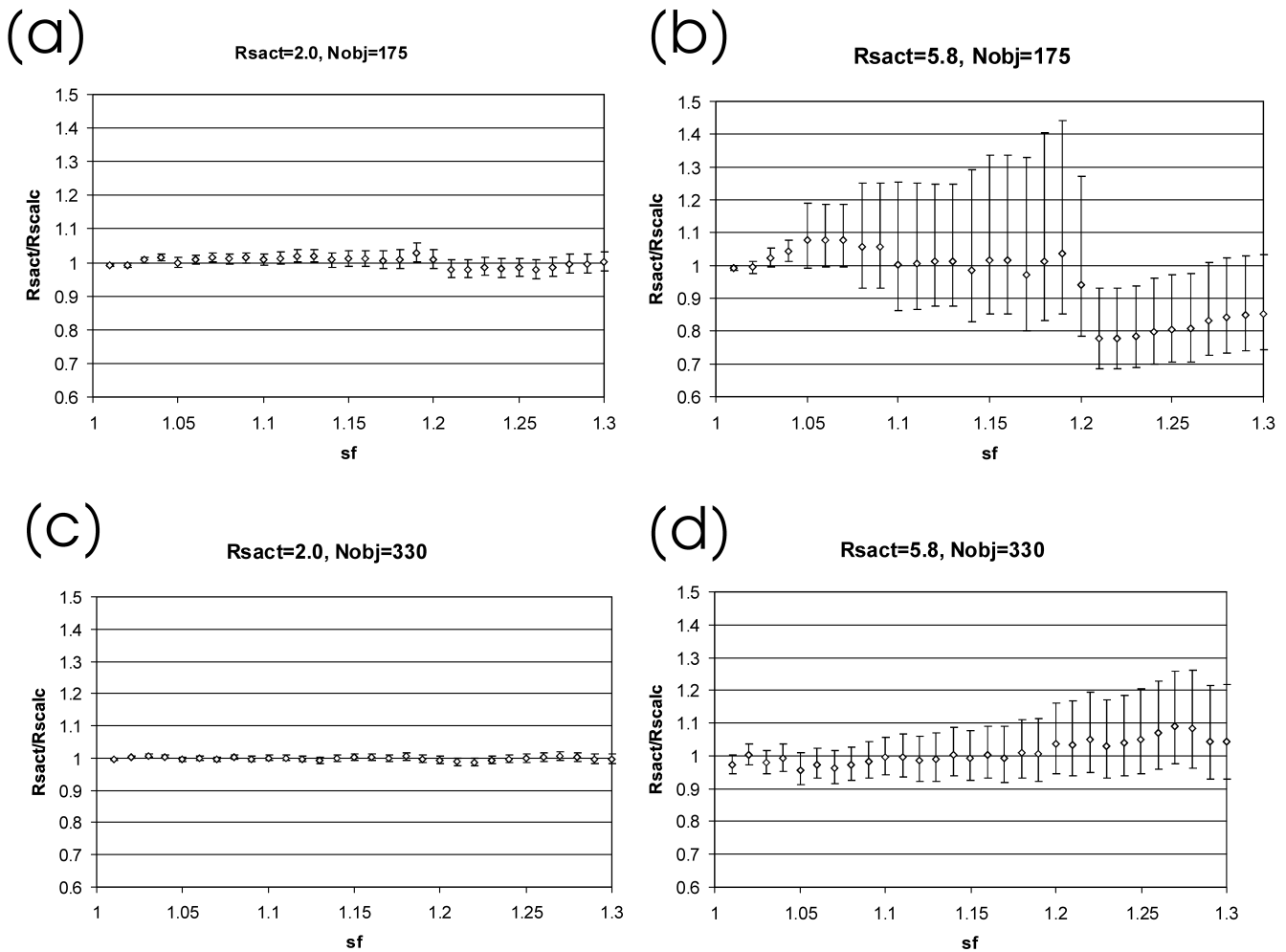


Fig. 10. Representative detailed results from the simulation study, showing R_{sact}/R_{scal} versus sf along with confidence intervals for: (a) $R_{sact} = 2.0$, $N_{obj} = 175$, (b) $R_{sact} = 5.8$, $N_{obj} = 175$, (c) $R_{sact} = 2.0$, $N_{obj} = 330$, and (d) $R_{sact} = 5.8$, $N_{obj} = 330$.

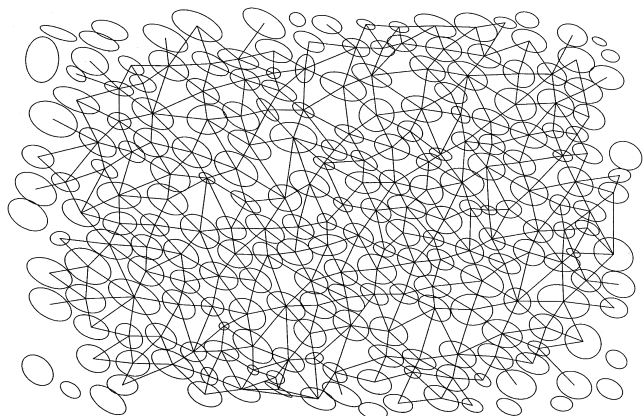


Fig. 11. Model of oolite data along with the corresponding modified Delaunay triangulation.

errors are calculated as:

$$R_{\text{serr}} = 100 \frac{|R_{\text{scalcu}} - R_{\text{scalcl}}|}{2R_{\text{scalc}}} \quad (11)$$

$$\phi_{\text{serr}} = 100 \frac{|\phi_{\text{scalcu}} - \phi_{\text{scalcl}}|}{2} \quad (12)$$

where R_{serr} is in percent and ϕ_{serr} is in degrees. From Fig. 8a it is clear that as n increases the error tends to decrease. Note that the number of objects analysed must be taken in the context of the size of the area analysed in so far as 1000 loosely packed objects measured from a large area may not be as effective as 200 closely packed objects from a smaller region. A packing density, defined as the ratio of object area to sample area, in conjunction with n may give a better indication of the numbers of objects required to achieve a desired level of accuracy. In the present simulation the packing densities are 34, 39, 46, 49, 52 and 54.5% corresponding to $n = 175, 205, 240, 270, 295$ and 330 , respectively.

Increasing the selection factor tends to produce greater errors in R_{scalc} and although an upward trend in the error for ϕ_{scalc} may also be distinguished, it is not as pronounced (see Fig. 8b). This indicates that for the most accurate results, the value of sf should be chosen close to one. As the tectonic strain increases (R_{sact}) the error associated with R_{scalc}

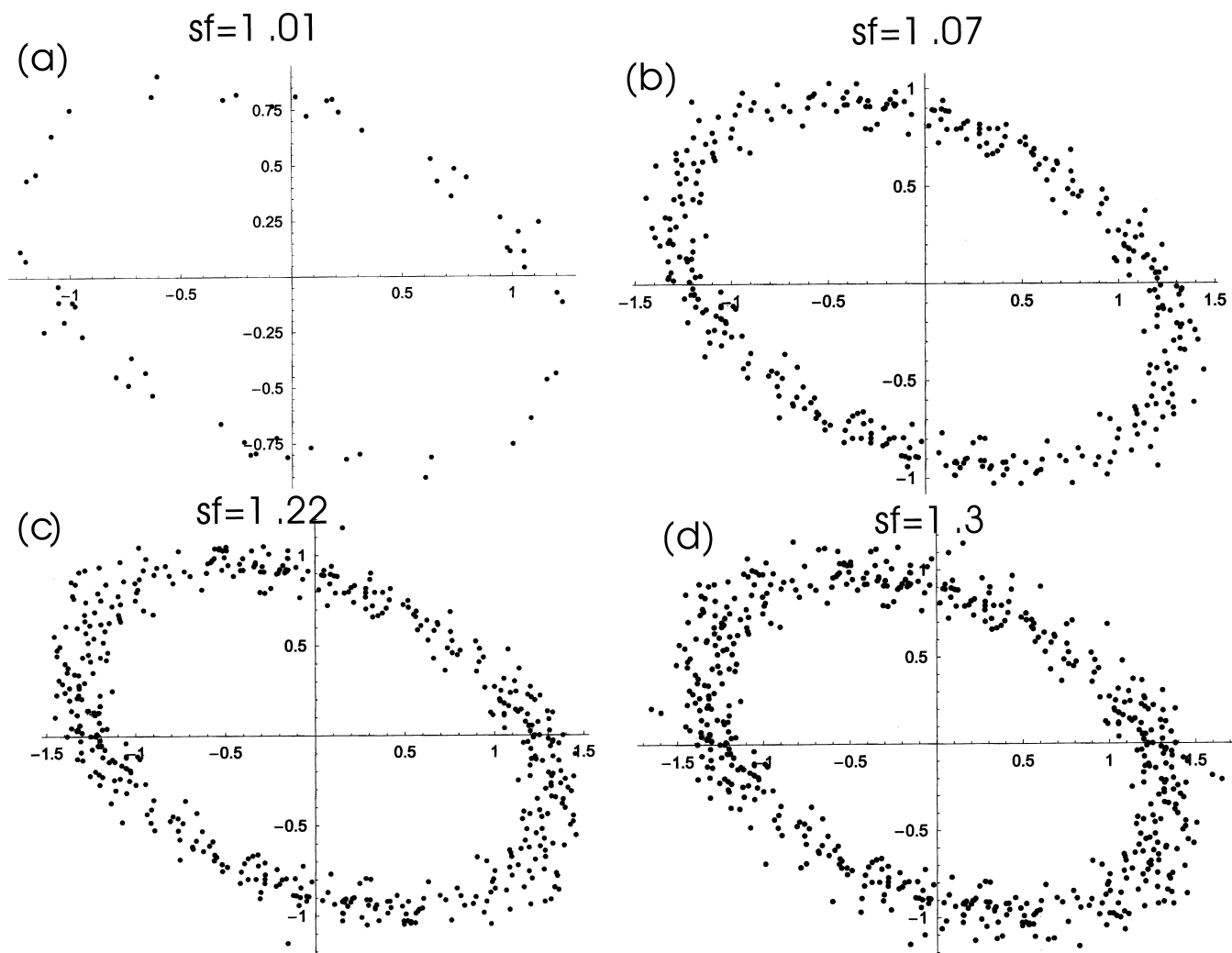


Fig. 12. Results of analysis for various values of sf .

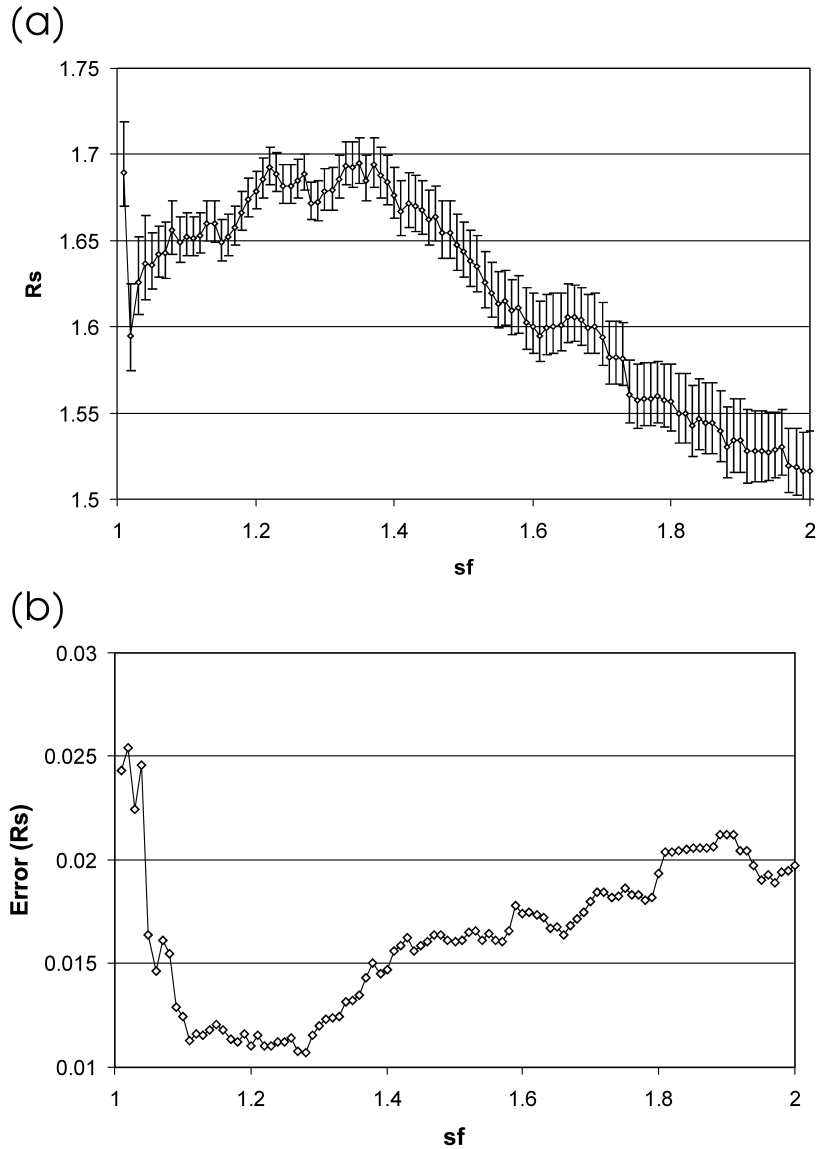


Fig. 13. (a) Variation of R_{scale} , with confidence intervals, for $sf = 1.01$ to 2. (b) Error variation with sf . Note that maximum R_{scale} broadly corresponds to the minimum in error.

increases correspondingly, e.g. from 1% ($R_{\text{sact}} \approx 1.2$) up to 10% ($R_{\text{sact}} \approx 5.8$) (see Fig. 8c). This is due in part to the variance of Delaunay triangulation (the changing nearest neighbour problem identified by Fry (1979)); however, it also incorporates statistical errors. This level of error is not unexpected because methods based on strain marker shape display errors of a comparable magnitude (Mulchrone and Meere, 2001; Mulchrone et al., 2002). On the basis of error characteristics alone, the DTNNM method (and presumably other Fry methods) out perform both the method of Robin (1977) and Mulchrone et al. (2002) using 100 strain markers at strain ratios less than approximately 5.8 because these methods are subject to a percentage error of approximately 13% for all strain ratios. By contrast with the error behaviour of strain ratio estimation, the error in estimating the orientation of the strain ellipse decreases rapidly as R_{sact} increases (Fig. 8c).

In Fig. 9 the percentage ratio of the number of data after enhancement (N_{proc}) to average number of data selected by the modified Delaunay triangulation (N_{del}) is compared with sf for $R_{\text{sact}} = 1, 2, 3, 4$ and 5. It is clear that as sf increases so too does the percentage of data available for analysis. In addition, as the imposed strain (R_{sact}) increases, the percentage of data available for analysis decreases.

These results suggest that as sf increases the additional data available for analysis tends to increase the error. However, at higher imposed strain less data is selected but the error is higher probably due to increased variability in the selected data. In summary, this simulation study suggests that the method proposed here is best suited to low to moderate strains and that the most accurate results are obtained by selecting low sf values.

Finally, representative detailed results are presented on $R_{\text{sact}}/R_{\text{scale}}$ versus sf plots (similar to those of Borradaile

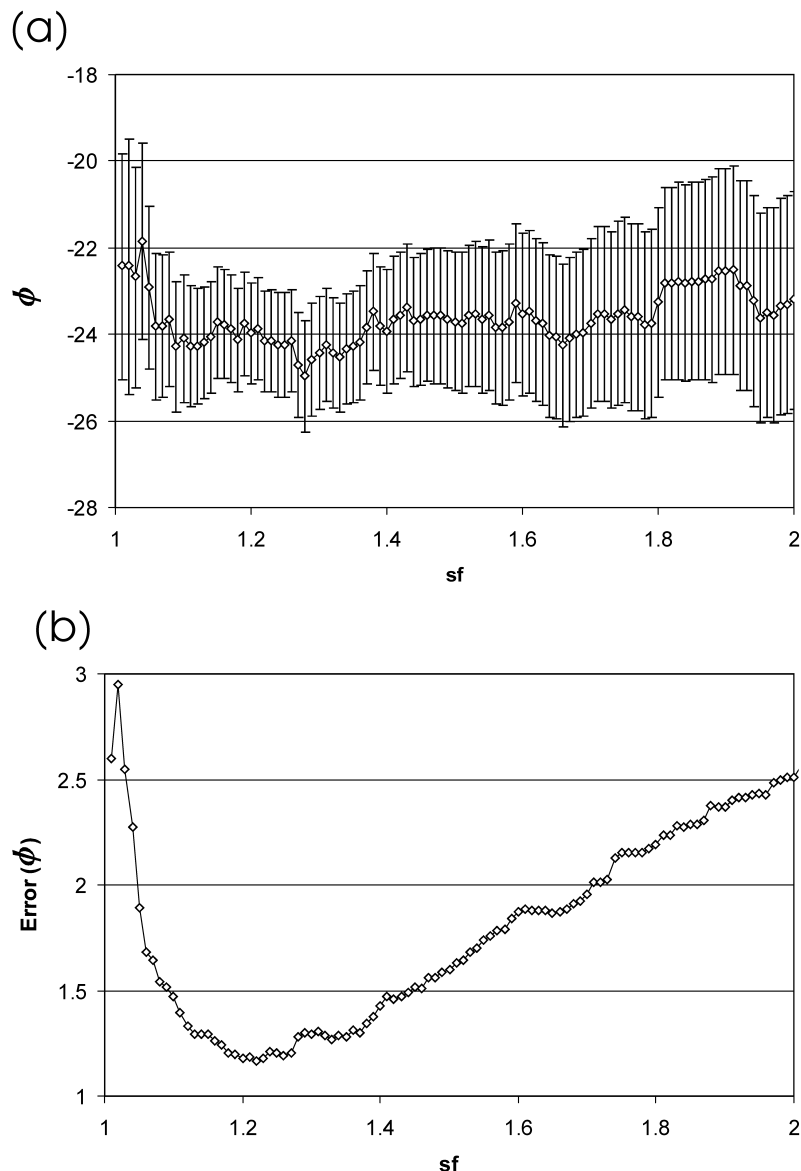


Fig. 14. (a) Variation of ϕ_{scale} , with confidence intervals, for $sf = 1.01$ to 2. (b) Error variation with sf . Note that maximum error in ϕ_{scale} corresponds to the minimum error in R_{scale} (Fig. 13b).

(1986); see Fig. 10). Values of one for $R_{\text{sact}}/R_{\text{scal}}$ indicates that the calculated strain matches the imposed strain exactly. Confidence intervals are also divided by R_{sact} and so represent percentage confidence intervals. These plots mirror the general trends outlined above. For $N_{\text{obj}} = 175$ (packing density of 34%), percentage confidence intervals are larger than those for $N_{\text{obj}} = 330$ (packing density of 54.5%) and in both cases the confidence intervals are much larger for higher imposed strain values, i.e. $R_{\text{sact}} = 2.0$ as compared with 5.8. However, in almost every case (except for the worst possible case of $N_{\text{obj}} = 175$ and $R_{\text{sact}} = 5.8$) the calculated strain is within 10% of the true value and the confidence intervals nearly always include $R_{\text{sact}}/R_{\text{scal}} = 1$. This simulation study strongly suggests that the method proposed here is valid and produces accurate results to a 95% level of confidence.

5. Application to real data

In order to demonstrate the applicability of the technique described here and also to investigate its properties with real data, the deformed ironstone oolite illustrated in fig. 7.7 of Ramsay and Huber (1983, p. 112) was analysed. The raw data parameters (a_i , b_i , x_i , y_i , ϕ_i) were calculated and the results of this are graphically illustrated in Fig. 11, along with the modified Delaunay triangulation. It is clear that nearest neighbours are connected in most cases. Plots of the data resulting from the analysis are illustrated in Fig. 12 for $sf = 1.01$, 1.07, 1.22 and 1.3. In each case a clearly defined ellipse is present. However, in order to decide which value of sf gives the best estimate of the strain ellipse the errors were studied in detail by varying sf in 0.01 steps.

By contrast with the simulation study above, analysis of

natural data shows a systematic variation with sf (see Figs. 13 and 14). In Fig. 13a, $R_{\text{scal}}c$ increases to a maximum at $sf \approx 1.3$ and then falls off for higher values of sf . According to the simulation study, the best estimate should be obtained by choosing a low value for sf ; however, in Fig. 13b, a systematic variation in the error is also observed which corresponds to the variation in $R_{\text{scal}}c$. Errors are much higher for low and high sf but are at a minimum for $1.1 < sf < 1.3$. Furthermore, a similar pattern is observed for $\phi_{\text{scal}}c$ (see Fig. 14a and b), although the value of $\phi_{\text{scal}}c$ is not as variable as $R_{\text{scal}}c$, and the minimum error is much more clearly defined. Based on this data it is suggested that the best fit strain ellipse may be estimated by choosing the value of sf that minimises the errors, that is $sf = 1.22$. The calculated strain values are therefore given by $R_{\text{scal}}c = 1.69 \pm 0.01$ and $\phi_{\text{scal}}c = -24.2 \pm 1.2^\circ$. The approach of error minimisation is not subjective and is a solution to the problem of how to select an appropriate value for the selection factor (McNaught, 1994). Furthermore, it is noted that Erslev and Ge (1990) report strain estimates for this sample that are somewhat lower (i.e. 1.598 up to 1.672); however, they used $sf = 1.01$ and 1.07 and had not minimised the errors.

It is interesting to note the differences between the behaviour of error in the natural example and the simulated data (see Fig. 8b). This is probably because the natural data is not perfectly anti-clustered and the oolites were not originally perfectly circular.

As a cross check on the above results, the shapes of the oolites were analysed using the methods of Robin (1977) and Mulchrone et al. (2002). The method of Robin (1977) gives $R_{\text{scal}}c = 1.68 \pm 0.05$ and $\phi_{\text{scal}}c = -23.5 \pm 1.0^\circ$, whereas the mean radial length approach of Mulchrone et al. (2002) gives $R_{\text{scal}}c = 1.69 \pm 0.05$ and $\phi_{\text{scal}}c = -23.6 \pm 1.1^\circ$. The strain ellipse calculated from the nearest neighbour analysis and the methods based on the shape of the oolites are identical within the precision of the methods. This result is somewhat surprising because the oolite analysed has undergone significant oolite-scale heterogeneous deformation by pressure solution. The nearest neighbour method measures strain due to matrix deformation and pressure solution; however, the methods based on oolite shape assume that the oolites deformed passively and homogeneously along with the matrix. This is clearly untrue. However, it appears that the pressure solution component of deformation has modified oolite shapes, such that they give an identical estimate of strain. This raises an interesting (and unanswered) question regarding the generality of this observation.

6. Conclusions

A method for objectively determining nearest neighbours, namely the Delaunay triangulation, is introduced in this paper. It is suggested that this approach makes the NNM of strain analysis a practical (and computationally more

efficient) alternative to the Fry and associated methods. Once nearest neighbours are selected, centre–centre distances can be processed by normalisation (Erslev, 1988) and enhancement (Erslev and Ge, 1990). The best fit ellipse is determined using a steepest gradient non-linear least squares algorithm applied to the polar equation on a centred ellipse. A simulation study indicates that the technique is a valid one and estimates the strain ellipse well at the 95% confidence interval. Application to a set of natural oolite data shows that there is a systematic variation of error with selection factor and it is suggested that the best estimate of the strain ellipse is obtained by choosing the selection factor which minimises the error. Moreover, the estimates of strain calculated with the DTNNM are comparable with estimates obtained using established methods.

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